Stochastic mechanics of graph rewriting

Nicolas Behr (LFCS, University of Edinburgh)

Joint work with Vincent Danos (ENS Paris) and Ilias Garnier (University of Edinburgh) July 9th 2016, LiCS'16, New York











Project supported by V. Danos' Advanced ERC grant RULE/320823

European Research Council Established by the European Commission

Concept

• TOPIC: continuous time Markov chains (CTMCs) for graph rewriting systems

• THE SPECIAL TWIST:

- link between non-deterministic transitions and associative unital algebras
- new insights into the combinatorics of graph rewriting, unification of the formalisms of physics and mathematical combinatorics with the computer science formalism
- powerful new formulae for fragmentation!
- **PLAN:** short motivation/introduction about **stochastic graph rewriting** (aka a particular class of CTMCs), then the new **rule algebraic formalism**

Graph rewriting and Markov chain theory

The basics of generic Markov chains

- Idea: consider a stochastic dynamical system
 with
 - a discrete state space (possibly ∞ many states though!)
 - the memoryless property, whence the transitions of the system can only depend on information available in the current state of the system (and not on the history of the state)
 - transitions are discrete jumps from one state to another – they happen instantaneuously, though at random times
- Closer analysis reveals that the jump times can then only be drawn from some **exponential jump time distribution** (that may be state-dependent)
- ⇒ characterization of a continuous time Markov chain!



The Master equation and other niceties

• Standard CTMC theory^[1]: one way to describe the dynamics is to give a **probability** distribution

$$|\Psi(t)
angle:=\sum_{\mathcal{S}\in\mathcal{S}}oldsymbol{
ho}_{\mathcal{S}}(t)|oldsymbol{S}
angle$$

of being in one of the discrete states (represented by basis vectors $|S\rangle$), and specifying the **Master equation** (aka **Schrödinger equation**)

$$\frac{d}{dt}|\Psi(t)\rangle = H|\Psi(t)\rangle,$$

where *H* is the **evolution operator**.

• How precisely *H* is determined for a given system is one of the main questions in general!



^[1] James R. Norris. Markov Chains. Cambridge Series in Statistical and Probabilistic Mathematics. Cambridge University Press, 1998

The "stochastic mechanics" viewpoint

Benefits:

∃ a full-blown formalism aka "**stochastic mechanics**"^[2] for studying CTMCs:

• **Observables** *O* are linear operators under which each pure state is an Eigenstate,

 $O|S\rangle = \omega_O(S)|S\rangle.$

• Expectation values of observables are computed by introducing the dual projection vector

$$\langle | \boldsymbol{S} \rangle := 1 \quad \forall \boldsymbol{S} \in \mathcal{S} \,,$$

such that for any state probability distribution $|\Psi(t)
angle$

$$\mathbb{E}_{|\Psi(t)\rangle}(O) \equiv \langle O \rangle(t) := \langle |O|\Psi(t) \rangle.$$

⇒ evolution of expectation values of observables via Master equation:

$$\frac{d}{dt}\langle O \rangle(t) = \langle OH \rangle(t)$$
.

• Additional property of the evolution operator *H*:

$$\langle | \boldsymbol{e}^{tH} | \Psi(t) \rangle \stackrel{!}{=} \mathbf{1} \quad \Rightarrow \quad \langle | \boldsymbol{H} = \mathbf{0} \,,$$

i.e. *H* preserves normalizations.

⇒ analogue of the Ehrenfest equation of quantum mechanics:

$$\frac{d}{dt}\langle O\rangle(t)=\langle [O,H]\rangle(t)\,,$$

where [A, B] := AB - BA is the **commutator**

^[2] John C Baez and Jacob Biamonte. "A course on quantum techniques for stochastic mechanics". In: s arXiv:1209.3632 (2012)

Specializing to graph rewriting systems

Basics:

- State space: G_≅ isomorphism classes of finite directed multi-graphs
- Discrete transitions: isomorphism classes

 of graph rewriting rules I → O, with
 input graph I ∈ G_≃, output graph O ∈ G_≃,
 and where r : I → O is an isomorphism
 class of a partial injective morphism of
 graphs
- Evolution operator *H*: to be dtermined!

Insight from statistical physics/chemistry:

- Promote the discrete state space to a vector space (usually over R) S, with one basis vector per possible discrete state.
- The transitions themselves should form an associative unital algebra, implementing the action of transitions on states via a (choice of) representation of this algebra , where a representation is an algebra homomorphism from the algebra of transitions *T* to the space of endomorphisms End(S) over the state vector space S:

$$\rho: T \to End(\mathcal{S})$$

such that

$$\rho(t_1 * t_2) = \rho(t_1)\rho(t_2).$$

The main objectives of this work

- Implement the full CTMC/ stochastic mechanics formalism for graph rewriting this will require in particular to find the state space, implementation of transitions via a transition algebra, defining the evolution operator, and finally the construction of the canonical representation of the transition algebra.
- Determine a description of the evolution of graph rewriting systems in this formalism.
- Study the potential benefits of this approach!

Stochastic mechanics formalism

Step 1: Directed multigraphs and the graph state space

- A finite directed multigraph G ≡ (V, E, s, t) is specified as a set of vertices V, a set of edges E, and source and target maps s : E → V andt : E → V.
- The graph state space G is defined as the span of the set of isomorphism classes of graphs G_≅,

 $\mathcal{G} := span_{\mathbb{R}}(G_{\cong})$.

Step 2: Normal rule diagrams

- A graph rewriting rule *I* ⇒ O consists of an input graph *I*, an output graph O, and an injective partial morphism *r* : *I* → O. We will always consider such rules up to isomorphism (of *I* and *O*, and the induced effect on *r*).
- Crucial idea: encode rules as normal rule diagrams draw *I* at the bottom, *O* at the top and *r* as dotted lines (with ×··· to indicate which vertices and edges are not mapped)!

An example of a rule and its associated normal rule diagram.

Step 3: Generic rule diagrams and diagram composition

- The composition of rule diagrams may be understood visually as "threading together" normal rule diagrams along diagram matches, linking output vertices of one diagram to input vertices of another diagram, and similarly for the edges, subject to some consistency conditions:
 - one-to-one any output vertex or edge is connected to at most one input vertex or edge
 - acyclicity viewed in total, the rule diagram must not contain any cycles at all
 - delayed edge morphism condition if an output edge is linked to an input edge of another diagram, its endpoint vertices must be linked, too, possibly via intermediate in-out vertex mappings
 - totality see the paper!



Example for a composition of three normal rule diagrams (yellow boxes) into a composite rule diagram.

Step 4: Normalization of rule diagrams

 Given a generic rule diagram *d*, its normalization ∂(*d*) is defined as follows:



Example for a composition of three normal rule diagrams into a composite rule diagram.

Step 4: Normalization of rule diagrams

- Given a generic rule diagram *d*, its normalization ∂(*d*) is defined as follows:
 - Determine the input interface I(d) and the output interface O(d) of d as those input and output graph parts that are not involved in the diagram matches of d.



Input interface $\mathcal{I}(d)$ and output interface $\mathcal{O}(d)$ of the composite rule diagram *d*.

Step 4: Normalization of rule diagrams

- Given a generic rule diagram *d*, its normalization ∂(*d*) is defined as follows:
 - Determine the input interface I(d) and the output interface O(d) of d as those input and output graph parts that are not involved in the diagram matches of d.
 - "Read out" the internal structure of ∂(d) via following the worldlines from input to output parts all complete paths from I(d) to O(d) will be part of the injective partial map r
 (d) of ∂(d) ≡ (I(d), O(d), r
 (d), Ø).
- Note that the normalized diagram ∂(d) may again be interpreted as a graph rewriting rule, i.e.

$$\partial(d) \stackrel{\widehat{}}{=} \overline{r}(d) : \mathcal{I}(d) \longrightarrow \mathcal{O}(d)$$
.



Normalization of the composite rule diagram.

Step 5: Rule diagram algebra

Rule diagram algebra

Denote by $\mathcal{D} := span_{\mathbb{R}}(D_{\cong})$ the vector space of isomorphism classes of rule diagrams D_{\cong} . Then one may define an **algebra composition** $*_{\mathcal{D}}$ as follows: given two rule diagrams $d_1, d_2 \in \mathcal{D}$, define

$$d_1 *_{\mathcal{D}} d_2 := \sum_{\text{all diagram matches } m} \text{iso-class of } d_1 \text{ composed with } d_2 \text{ along } m$$

This compositon is then extended by linearity to all of \mathcal{D} . We call $\mathcal{D} \equiv (\mathcal{D}, +, \cdot, *_{\mathcal{D}})$ the **rule diagram** algebra.

Theorem

The rule diagram algebra is a **noncommutative**, unital associative algebra, with unit the empty rule diagram d_{\emptyset} .

Rule algebra

Denote by $\mathcal{R} := span_{\mathbb{R}}(R_{\cong})$ the vector space of isomorphism classes of normal rule diagrams R_{\cong} . Then one may define an **algebra composition** $*_{\mathcal{R}}$ as follows: given two normal rule diagrams $r_1, r_2 \in \mathcal{R}$, define

$$\mathbf{r}_{1} \ast_{\mathcal{R}} \mathbf{r}_{2} := \partial(\psi(\mathbf{r}_{1}) \ast_{\mathcal{D}} \psi(\mathbf{r}_{2})),$$

where $\psi : \mathcal{R} \hookrightarrow \mathcal{D}$ is the natural inclusion. This compositon is then extended by linearity to all of \mathcal{R} . We call $\mathcal{R} \equiv (\mathcal{R}, +, \cdot, *_{\mathcal{R}})$ the **rule algebra**.

Theorem

The rule algebra is a **noncommutative**, unital associative algebra, with unit the empty rule diagram d_{\emptyset} .

Note: The proof is based on demonstrating that φ
 [¯] := ∂ ∘ (ψ ⊗ ψ) is an algebra homomorphism
 from the rule diagram algebra D to the rule algebra R.

Step 7: Canonical representation of the rule algebra

 Let |G> denote the basis vectors of the vector space G of isomorphism classes of finite directed multigraphs.

Canonical representation of the rule algebra

Taking inspiration from the **canonical representation** of the so-called **heisenberg-Weyl algebra**, one may easily define the **canonical representation of the rule algebra** as follows:

$$\rho(O \stackrel{r}{\leftarrow} I) | \varnothing \rangle := \begin{cases} |O\rangle, & \text{if } I = \varnothing \\ \mathbf{0} \cdot | \varnothing \rangle, & \text{else.} \end{cases}$$

Moreover, for $G \neq \emptyset$, we define

$$\rho(\boldsymbol{O} \stackrel{r}{\Leftarrow} \boldsymbol{I}) | \boldsymbol{G} \rangle := \rho((\boldsymbol{O} \stackrel{r}{\Leftarrow} \boldsymbol{I}) *_{\mathcal{R}} (\boldsymbol{G} \Leftarrow \varnothing)) | \varnothing \rangle.$$

Theorem

This definition yields a proper representation.

Step 8: Stochastic mechanics

Core definition: the evolution operator

Given a stochastic graph rewriting system *T* in the form of a set of graph rewriting rules $I \stackrel{r}{\Rightarrow} O$ and a set of base rates $\kappa_r \in \mathbb{R}_{\geq 0}$, we may define the evolution operrator *H* as follows:

$$H := \sum_{r \in \mathcal{T}} \kappa_r \left(\rho(r) - \rho(I \xleftarrow{id_{dom(r)}} I) \right) \,.$$

Theorem

H thus defined

- is an infinitesimal stochastic operator
- captures the dynamics of the system via the Master equation $\frac{d}{dt}|\Psi(t)\rangle = H|\Psi(t)\rangle$.

Consistency check

Specializing to **discrete graph rewriting**, we recover precisely the evolution operator of **chemical reaction systems**, as described in the **Doi-Peliti formalism**!^{[3][4][5]}

[3] Masao Doi. "Second quantization representation for classical many-particle system". In: Journal of Physics A: Mathematical and General 9.9 (1976), p. 1465

[4] Masao Doi. "Stochastic theory of diffusion-controlled reaction". In: Journal of Physics A: Mathematical and General 9.9 (1976), p. 1479

[5] L Peliti. "Path integral approach to birth-death processes on a lattice". In: Journal de Physique 46.9 (1985), pp. 1469–1483

Harvesting of results

Unification

• We identified a formalism in which a stochastic graph rewriting system is described in a fashion completely standard in both CTMC and statistical physics theory!

Fragmentation theorem rederivation

Lemma

For any rules $O_1 \stackrel{r_1}{\leftarrow} I$ and $O_2 \stackrel{r_2}{\leftarrow} I$ with $dom(r_1) = dom(r_2)$, we have that

$$\langle | \rho(O_1 \stackrel{r_1}{\Leftarrow} I) = \langle | \rho(O_2 \stackrel{r_2}{\Leftarrow} I) = \langle | \rho(I \stackrel{r_l}{\Leftarrow} I) ,$$

with $r_1 := id_{dom(r_1)} = id_{dom(r_2)}$.

Theorem: Jump-closure of observables (compare [6])

Let *H* be the evolution operator of a stochastic graph rewriting system. Then for any observable $O_l := I \stackrel{r_l}{\Leftarrow} I$ we have that

$$\langle |O_{I}H = \sum_{O_{I'} \in \mathcal{F}(O_{I})} \alpha_{O_{I},O_{I'},H} \langle |O_{i'},$$

for some constants $\alpha_{O_i,O_{l'},H} \in \mathbb{R}$, and with $\mathcal{F}(O_l)$ a finite family of observables.

Corollary: Fragmentation Theorem (compare [6])

$$\frac{d}{dt} \langle O_l \rangle(t) = \sum_{O_{l'} \in \mathcal{F}(O_l)} \alpha_{O_l, O_{l'}, H} \langle O_{i'} \rangle(t) \,.$$

[6]

New result: Generalized Fragmentation Theorem

 Any evolution operator H admits a unique decomposition H = Ĥ + H, with the non-observable part Ĥ and observable part H

$$\hat{H} = \rho(\hat{r}): \quad \hat{r} \in \mathcal{R} \backslash \mathcal{O}, \quad \hat{H} = \rho(\hat{r}): \quad \hat{r} \in \mathcal{O},$$

where $\mathcal{O} \subset \mathcal{R}$ denotes the subalgebra of **observables**.

Theorem: Generalized Fragmentation Theorem

For any evolution operator *H* and any observables $O_1, \ldots, O_n \in \mathcal{O}$, we have that

$$\frac{d}{dt}\langle O_1\ldots O_n\rangle(t) = \sum_{\sigma\in S_n} \frac{1}{n!} \sum_{m=1}^n \binom{n}{m} \langle C(\mathcal{O}_m^{\sigma}, \hat{H}) \prod_{i>m} O_{\sigma(i)}\rangle(t),$$

where

$$\mathcal{O}_m^{\sigma} := \{\mathcal{O}_{\sigma(i)}\}_{1 \leqslant i \leqslant m},$$

and with the nested commutators

$$C(\mathcal{O}_m^{sigma}, \hat{H}) := [\mathcal{O}_{\sigma(1)}, [\mathcal{O}_{\sigma(2)}, [\dots, [\mathcal{O}_{\sigma(m)}, \hat{H}] \dots]].$$

Conclusion and outlook

Summary

- We defined a full stochastic mechanics framework for graph rewriting (of DPO type).
- Based on the new formulae, we were able to achieve not only a standard formulation of the so-called **Fragmentation Theorem**, but more importantly found a compact expression for a **Generalized Fragmentation Theorem** for arbitrary moments of observables.
- The full formalism is an autonomous, mathematically well-posed framework for graph rewriting.

Outlook

- All results presented strictly speaking are only fully well-behaved for **finite** systems; while practical experience would suggest applicability also for infinite systems with suitable properties, this remains an open research question!
- Study mathematical properties of the rule diagram algebra more closely^[7].
- Generalization of the rule algebra idea to other types of rewriting (e.g. SPO types^[7], Kappa, biological tissue models...).
- The combinatorics of rewriting is in principle accessible along the lines of the analytical combinatorics formalism^[8].
- Generalization of the formalism to branching and fusing rewriting (and also a proper analysis of the category-theoretical foundations of graph rewriting) – project with V. Danos, I. Garnier and P. Sobocinski!

^[7] Nicolas Behr, Vincent Danos, Ilias Garnier, and Tobias Heindel. "The algebras of graph rewriting". In: (pprox Q4 2016)

^[8] Philippe Flajolet and Robert Sedgewick. Analytic combinatorics. Cambridge University Press, 2009

Thank you!

- John C Baez and Jacob Biamonte. "A course on quantum techniques for stochastic mechanics". In: *s arXiv:1209.3632* (2012).
- Nicolas Behr, Vincent Danos, and Ilias Garnier. "Stochastic mechanics of graph rewriting". In: *Thirty-First Annual ACM/IEEE Symposium on Logic in Computer Science (LICS)*. 2016.
- Nicolas Behr, Vincent Danos, Ilias Garnier, and Tobias Heindel. "The algebras of graph rewriting". In: (≈ Q4 2016).

- Vincent Danos et al. "Approximations for Stochastic Graph Rewriting". In: *Formal Methods and Software Engineering 16th International Conference on Formal Engineering Methods, ICFEM 2014.* Ed. by Stephan Merz and Jun Pang. Vol. 8829. Lecture Notes in Computer Science. Springer, 2014, pp. 1–10.
- Masao Doi. "Second quantization representation for classical many-particle system". In: *Journal of Physics A: Mathematical and General* 9.9 (1976), p. 1465.
- Masao Doi. "Stochastic theory of diffusion-controlled reaction". In: *Journal of Physics A: Mathematical and General* 9.9 (1976), p. 1479.
- Philippe Flajolet and Robert Sedgewick. *Analytic combinatorics*. Cambridge University Press, 2009.
- James R. Norris. *Markov Chains*. Cambridge Series in Statistical and Probabilistic Mathematics. Cambridge University Press, 1998.
- L Peliti. "Path integral approach to birth-death processes on a lattice". In: *Journal de Physique* 46.9 (1985), pp. 1469–1483.